

Chiral Torsion and its quantum effects

Amitabha Lahiri

S. N. Bose National Centre for Basic Sciences

Block JD, Sector III, Salt Lake, Kolkata, INDIA

Future Trends in Gravitational Physics

SNBNCBS, 10th February 2022



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e^a_μ relate spacetime with internal flat space



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e_μ^a relate spacetime with internal flat space
 - $\eta_{ab} e_\mu^a e_\nu^b = g_{\mu\nu}$; $g^{\mu\nu} e_\mu^a e_\nu^b = \eta^{ab}$,



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e^a_μ relate spacetime with internal flat space
 - $\eta_{ab} e^a_\mu e^b_\nu = g_{\mu\nu}$; $g^{\mu\nu} e^a_\mu e^b_\nu = \eta^{ab}$,
 - $e^\mu_a = \eta_{ab} g^{\mu\nu} e^b_\nu$ is inverse tetrad, $e^\mu_a e^a_\nu = \delta^\mu_\nu$



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e_μ^a relate spacetime with internal flat space
 - $\eta_{ab} e_\mu^a e_\nu^b = g_{\mu\nu}$; $g^{\mu\nu} e_\mu^a e_\nu^b = \eta^{ab}$,
 - $e_a^\mu = \eta_{ab} g^{\mu\nu} e_\nu^b$ is inverse tetrad, $e_a^\mu e_\nu^a = \delta_\nu^\mu$
- $\gamma_\mu = e_\mu^a \gamma_a$ satisfies $[\gamma_\mu, \gamma_\nu]_+ = 2g_{\mu\nu}$ (will use both γ^a and γ^μ as convenient)



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e_μ^a relate spacetime with internal flat space
 - $\eta_{ab} e_\mu^a e_\nu^b = g_{\mu\nu}$; $g^{\mu\nu} e_\mu^a e_\nu^b = \eta^{ab}$,
 - $e_a^\mu = \eta_{ab} g^{\mu\nu} e_\nu^b$ is inverse tetrad, $e_a^\mu e_\nu^a = \delta_\nu^\mu$
- $\gamma_\mu = e_\mu^a \gamma_a$ satisfies $[\gamma_\mu, \gamma_\nu]_+ = 2g_{\mu\nu}$ (will use both γ^a and γ^μ as convenient)
- Dirac operator $i\cancel{\partial}\psi$ on flat space becomes $i\cancel{D}\psi$ on curved spacetime



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e_μ^a relate spacetime with internal flat space
 - $\eta_{ab} e_\mu^a e_\nu^b = g_{\mu\nu}$; $g^{\mu\nu} e_\mu^a e_\nu^b = \eta^{ab}$,
 - $e_a^\mu = \eta_{ab} g^{\mu\nu} e_\nu^b$ is inverse tetrad, $e_a^\mu e_\nu^a = \delta_\nu^\mu$
- $\gamma_\mu = e_\mu^a \gamma_a$ satisfies $[\gamma_\mu, \gamma_\nu]_+ = 2g_{\mu\nu}$ (will use both γ^a and γ^μ as convenient)
- Dirac operator $i\cancel{\partial}\psi$ on flat space becomes $i\cancel{D}\psi$ on curved spacetime
- Requires spin connection A_μ : $D_\mu\psi = \partial_\mu\psi - \frac{i}{4} A_\mu^{ab} \sigma_{ab}\psi$
 $\sigma_{ab} = \frac{i}{2} [\gamma_a, \gamma_b]_-$



Fermions on curved spacetime

- Dirac matrices defined on "internal" flat space, $[\gamma^a, \gamma^b]_+ = \eta^{ab}$
- Tetrad fields e_μ^a relate spacetime with internal flat space
 - $\eta_{ab} e_\mu^a e_\nu^b = g_{\mu\nu}$; $g^{\mu\nu} e_\mu^a e_\nu^b = \eta^{ab}$,
 - $e_a^\mu = \eta_{ab} g^{\mu\nu} e_\nu^b$ is inverse tetrad, $e_a^\mu e_\nu^a = \delta_\nu^\mu$
- $\gamma_\mu = e_\mu^a \gamma_a$ satisfies $[\gamma_\mu, \gamma_\nu]_+ = 2g_{\mu\nu}$ (will use both γ^a and γ^μ as convenient)
- Dirac operator $i\cancel{\partial}\psi$ on flat space becomes $i\cancel{D}\psi$ on curved spacetime
- Requires spin connection A_μ : $D_\mu \psi = \partial_\mu \psi - \frac{i}{4} A_\mu^{ab} \sigma_{ab} \psi$
 $\sigma_{ab} = \frac{i}{2} [\gamma_a, \gamma_b]_-$
- $\nabla_\mu \gamma_\nu = 0 \Rightarrow \nabla_\mu e_\nu^a = 0$
 $\Rightarrow e_a^\lambda \partial_\mu e_\nu^a + A_{\mu b}^a e_\nu^b e_a^\lambda - \Gamma_{\mu\nu}^\lambda = 0$ "tetrad postulate"



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$
- Ricci scalar: $R(\Gamma) = F_{\mu\nu}^{ab} e_a^\mu e_b^\nu$



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$
- Ricci scalar: $R(\Gamma) = F_{\mu\nu}^{ab} e_a^\mu e_b^\nu$
- Action for gravity coupled to fermions:

$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \frac{i}{2} (\bar{\psi} \not{D} \psi - (\bar{\psi} \not{D} \psi)^\dagger) + im \bar{\psi} \psi \right]$$



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$
- Ricci scalar: $R(\Gamma) = F_{\mu\nu}^{ab} e_a^\mu e_b^\nu$
- Action for gravity coupled to fermions:

$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \frac{i}{2} (\bar{\psi} \not{D} \psi - (\bar{\psi} \not{D} \psi)^\dagger) + im \bar{\psi} \psi \right]$$

- metric signature $(-+++)$ [γ_0 is anti-Hermitian]



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$

- Ricci scalar: $R(\Gamma) = F_{\mu\nu}^{ab} e_a^\mu e_b^\nu$

- Action for gravity coupled to fermions:

$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \frac{i}{2} (\bar{\psi} \not{D} \psi - (\bar{\psi} \not{D} \psi)^\dagger) + im \bar{\psi} \psi \right]$$

- metric signature $(-+++)$ [γ_0 is anti-Hermitian]
- All fermions should be included (more about this later)



Einstein-Cartan-Kibble-Sciama formulation

- Can write gravity as a gauge theory with spin connection as the gauge field
- Field strength $F_{\mu\nu}^{ab} = \partial_\mu A_\nu^{ab} - \partial_\nu A_\mu^{ab} + A_\mu^{ac} A_\nu^{cb} - A_\nu^{ac} A_\mu^{cb}$
- Ricci scalar: $R(\Gamma) = F_{\mu\nu}^{ab} e_a^\mu e_b^\nu$
- Action for gravity coupled to fermions:

$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \frac{i}{2} (\bar{\psi} \not{D} \psi - (\bar{\psi} \not{D} \psi)^\dagger) + im \bar{\psi} \psi \right]$$

- metric signature $(-+++)$ [γ_0 is anti-Hermitian]
- All fermions should be included (more about this later)
- Work with one species for the moment



Equations

- Gauge theory action
$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$$



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$
- A_μ and e_μ are independent fields in this framework



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$
- A_μ and e_μ are independent fields in this framework
- $\delta A_\nu^{ab} :$ $A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$
- A_μ and e_μ are independent fields in this framework
- $\delta A_\nu^{ab} :$ $A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$
 - $\omega_\mu^{ab}[e]$ corresponds to the torsion-free Levi-Civita connection



Equations

- Gauge theory action
$$S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$$
- $$\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$$
- A_μ and e_μ are independent fields in this framework
- $\delta A_\nu^{ab} :$
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$$
 - $\omega_\mu^{ab}[e]$ corresponds to the torsion-free Levi-Civita connection
 - Christoffel symbols
$$\Gamma_{\mu\nu}^\sigma = e_a^\sigma \partial_\mu e_\nu^a + e_a^\sigma e_{\nu b} \omega_\mu^{ab}$$



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$
- A_μ and e_μ are independent fields in this framework
- $\delta A_\nu^{ab} :$ $A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$
 - $\omega_\mu^{ab}[e]$ corresponds to the torsion-free Levi-Civita connection
 - Christoffel symbols $\Gamma_{\mu\nu}^\sigma = e_a^\sigma \partial_\mu e_\nu^a + e_a^\sigma e_{\nu b} \omega_\mu^{ab}$
 - The fermion bilinear is torsion



Equations

- Gauge theory action $S = \int |e| d^4x \left[\frac{1}{2\kappa} F_{\mu\nu}^{ab} e_a^\mu e_b^\nu + \mathcal{L}_\psi \right]$
- $\mathcal{L}_\psi = \frac{i}{2} \left(\bar{\psi} \gamma^\mu \partial_\mu \psi - \partial_\mu \bar{\psi} \gamma^\mu \psi - \frac{i}{4} A_\mu^{ab} \bar{\psi} [\sigma_{ab}, \gamma_c]_+ \psi e^{\mu c} \right) + im \bar{\psi} \psi$
- A_μ and e_μ are independent fields in this framework
- $\delta A_\nu^{ab} :$ $A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$
 - $\omega_\mu^{ab}[e]$ corresponds to the torsion-free Levi-Civita connection
 - Christoffel symbols $\Gamma_{\mu\nu}^\sigma = e_a^\sigma \partial_\mu e_\nu^a + e_a^\sigma e_{\nu b} \omega_\mu^{ab}$
 - The fermion bilinear is torsion
 - Dynamically generated, completely antisymmetric (geodesic eq. unaffected)



Equations

- Varying e^a_μ gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$



Equations

- Varying e^a_μ gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$
 - Insert A^{ab}_μ into this $\Rightarrow T_{\mu\nu}$ has terms quartic in the fermion



Equations

- Varying e_μ^a gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$

- Insert A_μ^{ab} into this $\Rightarrow T_{\mu\nu}$ has terms quartic in the fermion

- In terms of the torsion-free Levi-Civita connection,

$$\hat{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\hat{R} = \kappa\hat{T}_{\mu\nu}(\psi, \bar{\psi}) - \frac{3\kappa^2}{16}g_{\mu\nu}\bar{\psi}\gamma_a\gamma_5\psi\bar{\psi}\gamma^a\gamma_5\psi$$



Equations

- Varying e_μ^a gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$

- Insert A_μ^{ab} into this $\Rightarrow T_{\mu\nu}$ has terms quartic in the fermion

- In terms of the torsion-free Levi-Civita connection,

$$\hat{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\hat{R} = \kappa\hat{T}_{\mu\nu}(\psi, \bar{\psi}) - \frac{3\kappa^2}{16}g_{\mu\nu}\bar{\psi}\gamma_a\gamma_5\psi\bar{\psi}\gamma^a\gamma_5\psi$$

- Last term comes from using $[\gamma_c, \sigma_{ab}]_+ = 2\epsilon_{abcd}\gamma^d\gamma^5$



Equations

- Varying e_μ^a gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$
 - Insert A_μ^{ab} into this $\Rightarrow T_{\mu\nu}$ has terms quartic in the fermion
 - In terms of the torsion-free Levi-Civita connection,

$$\hat{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\hat{R} = \kappa\hat{T}_{\mu\nu}(\psi, \bar{\psi}) - \frac{3\kappa^2}{16}g_{\mu\nu}\bar{\psi}\gamma_a\gamma_5\psi\bar{\psi}\gamma^a\gamma_5\psi$$
 - Last term comes from using $[\gamma_c, \sigma_{ab}]_+ = 2\epsilon_{abcd}\gamma^d\gamma^5$
- Dirac Eq.: $\not{\partial}\psi - \frac{i}{4}\omega_\mu^{ab}\gamma_\mu\sigma_{ab}\psi + m\psi + \frac{3i\kappa}{8}(\bar{\psi}\gamma^a\gamma^5\psi)\gamma_a\gamma^5\psi = 0$



Equations

- Varying e_μ^a gives Einstein Eq. : $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \kappa T_{\mu\nu}(\psi, \bar{\psi})$
 - Insert A_μ^{ab} into this $\Rightarrow T_{\mu\nu}$ has terms quartic in the fermion
 - In terms of the torsion-free Levi-Civita connection,

$$\hat{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\hat{R} = \kappa\hat{T}_{\mu\nu}(\psi, \bar{\psi}) - \frac{3\kappa^2}{16}g_{\mu\nu}\bar{\psi}\gamma_a\gamma_5\psi\bar{\psi}\gamma^a\gamma_5\psi$$
 - Last term comes from using $[\gamma_c, \sigma_{ab}]_+ = 2\epsilon_{abcd}\gamma^d\gamma^5$
- Dirac Eq.: $\not{\partial}\psi - \frac{i}{4}\omega_\mu^{ab}\gamma_\mu\sigma_{ab}\psi + m\psi + \frac{3i\kappa}{8}(\bar{\psi}\gamma^a\gamma^5\psi)\gamma_a\gamma^5\psi = 0$
- Dirac Eq. in curved spacetime is naturally nonlinear (old stuff)



Many fermions

- All fermions must appear in the action



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe

- Then
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$$



Many fermions

- All fermions must appear in the action
- $\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$
- The sum is over all species of fermions in the universe
- Then $A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$
- Dirac Eq. is modified in the same way, for every type of fermion



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe

- Then
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$$

- Dirac Eq. is modified in the same way, for every type of fermion

- $$\not{\partial} \psi - \frac{i}{4} \omega_\mu^{ab} \gamma^\mu \sigma_{ab} \psi + m \psi + \frac{3i\kappa}{8} \left(\sum_f \bar{f} \gamma^a \gamma^5 f \right) \gamma_a \gamma^5 \psi = 0$$



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe

- Then
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$$

- Dirac Eq. is modified in the same way, for every type of fermion

- $$\not{\partial} \psi - \frac{i}{4} \omega_\mu^{ab} \gamma^\mu \sigma_{ab} \psi + m \psi + \frac{3i\kappa}{8} \left(\sum_f \bar{f} \gamma^a \gamma^5 f \right) \gamma_a \gamma^5 \psi = 0$$

- Term in brackets can be approximated by background average (for $f \neq \psi$)



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe

- Then
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$$

- Dirac Eq. is modified in the same way, for every type of fermion

- $$\not{\partial} \psi - \frac{i}{4} \omega_\mu^{ab} \gamma^\mu \sigma_{ab} \psi + m \psi + \frac{3i\kappa}{8} \left(\sum_f \bar{f} \gamma^a \gamma^5 f \right) \gamma_a \gamma^5 \psi = 0$$

- Term in brackets can be approximated by background average (for $f \neq \psi$)

- But $\langle \sum_f \bar{f} \gamma^a \gamma^5 f \rangle$ is spin density ~ 0 for generic matter distribution



Many fermions

- All fermions must appear in the action

- $$\mathcal{L}_\psi \rightarrow \frac{i}{2} \sum_f \left(\bar{f} \gamma^\mu \partial_\mu f - \partial_\mu \bar{f} \gamma^\mu f - \frac{i}{4} A_\mu^{ab} \bar{f} [\sigma_{ab}, \gamma_c]_+ f e^{\mu c} + 2m \bar{f} f \right)$$

- The sum is over all species of fermions in the universe

- Then
$$A_\mu^{ab} = \omega_\mu^{ab}[e] + \frac{\kappa}{8} e_\mu^c \sum_f \bar{f} [\gamma_c, \sigma^{ab}]_+ f$$

- Dirac Eq. is modified in the same way, for every type of fermion

- $$\not{\partial} \psi - \frac{i}{4} \omega_\mu^{ab} \gamma^\mu \sigma_{ab} \psi + m \psi + \frac{3i\kappa}{8} \left(\sum_f \bar{f} \gamma^a \gamma^5 f \right) \gamma_a \gamma^5 \psi = 0$$

- Term in brackets can be approximated by background average (for $f \neq \psi$)

- But $\langle \sum_f \bar{f} \gamma^a \gamma^5 f \rangle$ is spin density ~ 0 for generic matter distribution

- However, something else should be taken into account



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]})^{ab} + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]})^{ab} + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$
- Λ appears in action as $\frac{1}{2\kappa} \eta_{cd} e_a^{\mu} e_b^{\nu} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db} + \frac{1}{8} \sum_f e_c^{\mu} \Lambda_{\mu}^{ab} \bar{\psi}_f [\gamma^c, \sigma_{ab}]_+ \psi_f$



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]}^{ab})^{ab} + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$
- Λ appears in action as $\frac{1}{2\kappa} \eta_{cd} e_a^{\mu} e_b^{\nu} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db} + \frac{1}{8} \sum_f e_c^{\mu} \Lambda_{\mu}^{ab} \bar{\psi}_f [\gamma^c, \sigma_{ab}]_+ \psi_f$
- Purely algebraic eq. for Λ (non-propagating auxiliary field)



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]}^{ab}) + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$
- Λ appears in action as $\frac{1}{2\kappa} \eta_{cd} e_a^{\mu} e_b^{\nu} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db} + \frac{1}{8} \sum_f e_c^{\mu} \Lambda_{\mu}^{ab} \bar{\psi}_f [\gamma^c, \sigma^{ab}]_+ \psi_f$
- Purely algebraic eq. for Λ (non-propagating auxiliary field)
- Solution is $\Lambda_{\mu}^{ab} = \frac{\kappa}{8} e_{\mu}^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$, which can be put back in the action



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]}^{ab}) + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$
- Λ appears in action as $\frac{1}{2\kappa} \eta_{cd} e_a^{\mu} e_b^{\nu} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db} + \frac{1}{8} \sum_f e_c^{\mu} \Lambda_{\mu}^{ab} \bar{\psi}_f [\gamma^c, \sigma_{ab}]_+ \psi_f$
- Purely algebraic eq. for Λ (non-propagating auxiliary field)
- Solution is $\Lambda_{\mu}^{ab} = \frac{\kappa}{8} e_{\mu}^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$, which can be put back in the action
- Equivalent to starting with torsion free connection and a four-fermion interaction term



Chiral Torsion

- First go back and write $A_{\mu}^{ab} = \omega_{\mu}^{ab}(e) + \Lambda_{\mu}^{ab}$ (contorsion)
- Then the field strength is $F_{\mu\nu}^{ab}(A) = F_{\mu\nu}^{ab}(\omega) + (D_{[\mu}^{\omega} \Lambda_{\nu]}^{ab}) + \eta_{cd} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db}$
- Λ appears in action as $\frac{1}{2\kappa} \eta_{cd} e_a^{\mu} e_b^{\nu} \Lambda_{[\mu}^{ac} \Lambda_{\nu]}^{db} + \frac{1}{8} \sum_f e_c^{\mu} \Lambda_{\mu}^{ab} \bar{\psi}_f [\gamma^c, \sigma_{ab}]_+ \psi_f$
- Purely algebraic eq. for Λ (non-propagating auxiliary field)
- Solution is $\Lambda_{\mu}^{ab} = \frac{\kappa}{8} e_{\mu}^c \bar{\psi} [\gamma_c, \sigma^{ab}]_+ \psi$, which can be put back in the action
- Equivalent to starting with torsion free connection and a four-fermion interaction term
- All the equations are exactly the same as before; so what did we gain?



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own
- No reason to assume Λ couples identically to all fermions



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own
- No reason to assume Λ couples identically to all fermions
- Or to left- and right-handed fermions



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own
- No reason to assume Λ couples identically to all fermions
- Or to left- and right-handed fermions
- Generic fermion-contorsion coupling:

$$\frac{1}{8} \sum_f \Lambda_\mu^{ab} e_c^\mu \left(\lambda_{fL} \bar{f}_L [\gamma^c, \sigma_{ab}]_+ f_L + \lambda_{fR} \bar{f}_R [\gamma^c, \sigma_{ab}]_+ f_R \right)$$



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own
- No reason to assume Λ couples identically to all fermions
- Or to left- and right-handed fermions

- Generic fermion-contorsion coupling:

$$\frac{1}{8} \sum_f \Lambda_\mu^{ab} e_c^\mu \left(\lambda_{fL} \bar{f}_L [\gamma^c, \sigma_{ab}]_+ f_L + \lambda_{fR} \bar{f}_R [\gamma^c, \sigma_{ab}]_+ f_R \right)$$

- Torsion is totally antisymmetric $\sim \epsilon_{abcd} \sum_f (\lambda_{fL} \bar{f}_L \gamma^d \gamma^5 f_L + \lambda_{fR} \bar{f}_R \gamma^d \gamma^5 f_R)$



Chiral Torsion

- (Con)torsion couples to fermions as a vector field independently of gravity
- Terms containing Λ are invariant on their own
- No reason to assume Λ couples identically to all fermions
- Or to left- and right-handed fermions

- Generic fermion-contorsion coupling:

$$\frac{1}{8} \sum_f \Lambda_{\mu}^{ab} e_c^{\mu} \left(\lambda_{fL} \bar{f}_L [\gamma^c, \sigma_{ab}]_+ f_L + \lambda_{fR} \bar{f}_R [\gamma^c, \sigma_{ab}]_+ f_R \right)$$

- Torsion is totally antisymmetric $\sim \epsilon_{abcd} \sum_f (\lambda_{fL} \bar{f}_L \gamma^d \gamma^5 f_L + \lambda_{fR} \bar{f}_R \gamma^d \gamma^5 f_R)$
- So geodesic equation is unaffected – all particles fall at the same rate



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_5 f + \lambda_{fR} \bar{f}_R \gamma_5 f) \right)^2$$



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_5 f + \lambda_{fR} \bar{f}_R \gamma_5 f) \right)^2$$

- This term must always be present, and include all fermions



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_5 f + \lambda_{fR} \bar{f}_R \gamma_5 f) \right)^2$$

- This term must always be present, and include all fermions
- Unless $\lambda_{fL}, \lambda_{fR}$ are set to zero by hand



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_a \gamma^5 f + \lambda_{fR} \bar{f}_R \gamma_a \gamma^5 f_R) \right)^2$$

- This term must always be present, and include all fermions
- Unless $\lambda_{fL}, \lambda_{fR}$ are set to zero by hand
- Affects dispersion relations: e.g. neutrinos in matter



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_a \gamma^5 f + \lambda_{fR} \bar{f}_R \gamma_a \gamma^5 f_R) \right)^2$$

- This term must always be present, and include all fermions
- Unless $\lambda_{fL}, \lambda_{fR}$ are set to zero by hand
- Affects dispersion relations: e.g. neutrinos in matter
- Assume only left-handed neutrinos, then forward scattering amplitude in normal matter

$$\mathcal{M} = -(\bar{\nu}_\alpha \gamma_a \nu_\alpha) \lambda_{\nu_\alpha} \left\langle \sum_{e,p,n} (\lambda_{fN} \bar{\psi}_f \gamma^a \psi_f - \lambda_{fA} \bar{\psi}_f \gamma^a \gamma^5 \psi_f) \right\rangle$$



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_a \gamma^5 f + \lambda_{fR} \bar{f}_R \gamma_a \gamma^5 f_R) \right)^2$$

- This term must always be present, and include all fermions
- Unless $\lambda_{fL}, \lambda_{fR}$ are set to zero by hand
- Affects dispersion relations: e.g. neutrinos in matter
- Assume only left-handed neutrinos, then forward scattering amplitude in normal matter

$$\mathcal{M} = -(\bar{\nu}_\alpha \gamma_a \nu_\alpha) \lambda_{\nu_\alpha} \left\langle \sum_{e,p,n} (\lambda_{fV} \bar{\psi}_f \gamma^a \psi_f - \lambda_{fA} \bar{\psi}_f \gamma^a \gamma^5 \psi_f) \right\rangle$$

- Second term vanishes in typical matter distributions (axial charge and spin density)



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_a \gamma^5 f + \lambda_{fR} \bar{f}_R \gamma_a \gamma^5 f_R) \right)^2$$

- This term must always be present, and include all fermions
- Unless $\lambda_{fL}, \lambda_{fR}$ are set to zero by hand
- Affects dispersion relations: e.g. neutrinos in matter
- Assume only left-handed neutrinos, then forward scattering amplitude in normal matter

$$\mathcal{M} = -(\bar{\nu}_\alpha \gamma_a \nu_\alpha) \lambda_{\nu_\alpha} \left\langle \sum_{e,p,n} (\lambda_{fV} \bar{\psi}_f \gamma^a \psi_f - \lambda_{fA} \bar{\psi}_f \gamma^a \gamma^5 \psi_f) \right\rangle$$

- Second term vanishes in typical matter distributions (axial charge and spin density)
- First term averages to density, $\delta \mathcal{H}_{\text{eff}} = \left(\sum_{e,p,n} \lambda_f n_f \right) \lambda_{\nu_\alpha} \nu_\alpha^\dagger \nu_\alpha$



Applications

- Effective four-fermion interaction

$$-\frac{1}{8} \left(\sum_f (\lambda_{fL} \bar{f}_L \gamma_a \gamma^5 f + \lambda_{fR} \bar{f}_R \gamma_a \gamma^5 f_R) \right)^2$$

- This term must always be present, and include all fermions
- Unless λ_{fL} , λ_{fR} are set to zero by hand
- Affects dispersion relations: e.g. neutrinos in matter
- Assume only left-handed neutrinos, then forward scattering amplitude in normal matter

$$\mathcal{M} = -(\bar{\nu}_\alpha \gamma_a \nu_\alpha) \lambda_{\nu_\alpha} \left\langle \sum_{e,p,n} (\lambda_{fV} \bar{\psi}_f \gamma^a \psi_f - \lambda_{fA} \bar{\psi}_f \gamma^a \gamma^5 \psi_f) \right\rangle$$

- Second term vanishes in typical matter distributions (axial charge and spin density)

- First term averages to density, $\delta \mathcal{H}_{\text{eff}} = \left(\sum_{e,p,n} \lambda_f n_f \right) \lambda_{\nu_\alpha} \nu_\alpha^\dagger \nu_\alpha$

- Contributes to neutrino oscillations



Applications and speculations

- Treat as four-fermion interaction in flat space



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation
- Renormalization issues linked to QG



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation
- Renormalization issues linked to QG
- Main issue: How small is this interaction?



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation
- Renormalization issues linked to QG
- Main issue: How small is this interaction?
- How large does it need to be? (Ans. Not very.)



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation
- Renormalization issues linked to QG
- Main issue: How small is this interaction?
- How large does it need to be? (Ans. Not very.)
- λ_f couplings do not come from a quantum theory, unlike G_F



Applications and speculations

- Treat as four-fermion interaction in flat space
- Contribution to inertia and dispersion relations in matter
- Neutrino oscillations, stellar models, early universe
- Mass of dense stars may be different from calculated (models of dark matter?)
- Parity violation
- Renormalization issues linked to QG
- Main issue: How small is this interaction?
- How large does it need to be? (Ans. Not very.)
- λ_f couplings do not come from a quantum theory, unlike G_F
- Can be set only by comparison with experimental data



Work done with S. Chakrabarty: 1907.02341, 1904.06036
Also I. Ghose, R. Barik, A. Chakraborty (In preparation)

Thank You

