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Protein-cofactor binding and ultrafast electron transfer in riboflavin binding protein under the spatial confinement of nanoscopic reverse micelles

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In this contribution, we study the effect of confinement on the ultrafast electron transfer (ET) dynamics of riboflavin binding protein (RBP) to the bound cofactor riboflavin (Rf, vitamin B2), an important metabolic process, in anionic sodium bis(2-ethylhexyl) sulfosuccinate reverse micelles (AOT-RMs) of various hydration levels. Notably, in addition to excluded volume effect, various nonspecific interactions like ionic charge of the confining surface can influence the biochemical reactions in the confined environment of the cell. To this end, we have also studied the ET dynamics of RBP-Rf complex under the confinement of a cationic hexadecyltrimethylammonium bromide (CTAB) RMs with similar water pool size to the anionic AOT-RMs towards simulating equal restricted volume effect. It has been found that the spatial confinement of RBP in the AOT-RM of $w_0 = 10$ leads to the loss of its tertiary structure and hence vitamin binding capacity. Although, RBP regains its binding capacity and tertiary structure in AOT-RMs of w_0 \geq 20 due to its complete hydration, the ultrafast ET from RBP to Rf merely occurs in such systems. However, to our surprise, the ET process is found to occur in cationic CTAB-RMs of similar volume restriction. It is found that under the spatial confinement of anionic AOT-RM, the isoalloxazine ring of Rf is improperly placed in the protein nanospace so that ET between RBP and Rf is not permitted. This anomaly in the binding behaviour of Rf to RBP in AOT-RMs is believed to be the influence of repulsive potential of the anionic AOT-RM surface to the protein. Our finding thus suggests that under similar size restriction, both the hydration and surface charge of the confining volume could have major implication in the intraprotein ET dynamics in real cellular environments. Copyright © 2013 John Wiley & Sons, Ltd.

Keywords: riboflavin; riboflavin binding protein; electron transfer dynamics; confinement; hydration; surface charge of confinement; reverse micelles

INTRODUCTION

At the cellular interior, biopolymers such as proteins and DNA carry out various biological functions in a small space that can be approximated by nanosized confinement (or nanocavity). For example, DNA packs in virus capsids (Simpson *et al.*, 2000) and orients in pores (Mathé et al., 2005), proteins fold in chaperonin cages (Brinker et al., 2001; Ellis and Minton, 2003) and ribosomal exit tunnels (Nissen et al., 2000). Altogether, reactions in such in vivo confined environments differ from those processes in a simple aqueous solution, and with recent advances in various computational/experimental methods and resources: investigations of macromolecular crowding and confinement effects on protein conformational changes (Cheung and Thirumalai, 2007; Homouz et al., 2009), folding (Munishkina et al., 2004; Ping et al., 2004; Zhou, 2004; Zhang and Cheung, 2007; Qin and Zhou, 2009), thermodynamics and kinetics of protein folding under confinement (Mittal and Best, 2008), its association (Minton, 1993; Griffin et al., 2005; Wang et al., 2009) and dynamics (Bernadó et al., 2004; McGuffee and Elcock, 2006) have deepened our understanding of biopolymer dynamics. Although considerable progress has been made in this direction of protein folding reaction, little is known about biochemical reaction dynamics like ultrafast electron transfer (ET) under cell-like confinement. ET is crucial to life and is ubiquitous in enzymatic catalysis (Bendall, 1996; Stubbe and van der Donk, 1998; Balzani, 2001), especially in enzymes with redox reactions (Sinnott, 1998). Flavoproteins with flavin chromophores are examples of such enzymes and are involved in various catalytic processes (Müller, 1991; Ghisla *et al.*, 1999). The understanding of ET reaction dynamics of flavins in proteins and their redox reactions in cell-like confinement is crucial to the enzyme function.

In this contribution, we report the picosecond-resolved studies of ET dynamics of riboflavin (Rf; vitamin B2) in Rf-binding protein (RBP, a flavoprotein) under the confinement of sodium *bis*(2-ethylhexyl) sulfosuccinate reverse micelles (AOT-RMs) at various hydration levels ($w_0 = [water]/[surfactant]$). Encapsulation limits the available solvent, and the ability to precisely control micellar water

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Department of Chemical, Biological and Macromolecular Sciences, S.N. Bose National Centre for Basic Sciences, Block JD, Sector III Salt Lake, Kolkata 700098, India pool radius through hydration provides the flexibility to probe the effects of confinement in a systematic way. RM encapsulation is thus proposed to be an ideal model for crowded cellular confinement studies (Van Horn et al., 2009; Tian and Garcia, 2011; Yeung and Axelsen, 2012). Recently, it has been suggested that the dominant factors influencing protein behaviour in vivo are a combination of excluded volume effects and weak attractive forces (like hydrodynamic interactions and electrostatics) (Minton, 1983; Crowley et al., 2008; Li and Pielak, 2009; Jiao et al., 2010; Feig and Sugita, 2011). Nevertheless, how different electrostatic charge of the confining volumes affects proteincofactor binding and the associated reaction dynamics like ultrafast ET has not been studied so far. To this end, we also monitor the ET dynamics of Rf in RBP under the confinement of a cationic hexadecyltrimethylammonium bromide (CTAB) RMs with similar water pool size to the anionic AOT-RMs towards simulating equal restricted volume effect. We have monitored the secondary and tertiary structures and the vitamin binding capacity of RBP at different hydration levels (w_0) of the RMs using various spectroscopic techniques like circular dichroism, UV-Vis absorption, steady-state and picosecond-resolved fluorescence studies. The different dynamics of ET observed in such environments has been correlated with the hydration and structure of the protein in the corresponding nanoconfining environments.

MATERIALS AND METHODS

RBP (Apo form) from chicken egg white (lyophilized powder) was purchased from Sigma. Rf (Sigma), CTAB (Fluka), isooctane (Spectrochem, 99.5%), sodium bis(2-ethylhexyl) sulfosuccinate salt (AOT; Fluka 99%), disodium hydrogen phosphate dehydrate (Sigma 99%) and sodium dihydrogen phosphate dehydrate (Sigma 99%) were used as received. Aqueous stock solutions of RBP were prepared in a phosphate buffer (10 mM) at pH 7.0 using double distilled water. The concentration of RBP in buffer was determined using the extinction coefficient value of 49000 M^{-1} cm⁻¹ at 280 nm (Duyvis *et al.*, 2002). Rf concentration was calculated from its absorbance using the extinction coefficient value of $12200 \text{ M}^{-1} \text{ cm}^{-1}$ at 450 nm (Duyvis *et al.*, 2002). Unless otherwise mentioned, we have used $15 \,\mu M$ RBP and 7.5 µM Rf solutions for all the spectroscopic studies. Reverse micellar (RM) solutions were prepared by adding requisite volumes of aqueous solution of RBP or RBP-Rf complex into a given volume of surfactant solution (100 mM) in isooctane with gentle stirring to achieve RMs with required degree of hydration, w_0 ($w_0 = [water]/$ [surfactant]). For CTAB-RM preparation, hexanol was additionally added as a cosurfactant in 1:9 ratios (v/v) to the isooctane solution.

Far-, near- and visible-UV circular dichroism (CD) measurements were performed on a JASCO 815 spectrometer. Far-UV CD studies were measured between 200 and 260 nm wavelength with protein concentration of 0.15 mg ml^{-1} using a cell of 0.1 cmpath length. For the far-UV measurement, we used 50 mM AOT solution to avoid the noise in the CD spectrum due to the absorbance of RM in that region. Both near-UV and visible-UV CD measurements were made in 1.0 cm path length cell. Absorption and emission spectra were recorded with a Shimadzu UV-2450 spectrophotometer and a JobinYvon Fluoromax-3 fluorimeter, respectively. We used a commercially available picosecond diode laser-pumped (LifeSpec-ps) time-resolved fluorescence spectrophotometer from Edinburgh Instruments, UK, for timeresolved measurement. For 445 nm excitation, a diode laser from Edinburgh Instruments, UK, was used with instrument response function (IRF) of 80 ps. For RBP, we used a femtosecond-coupled time-correlated single-photon counting setup in which the sample was excited by the third harmonic laser beam (300 nm) of the 900 nm (0.5 nJ per pulse) using a mode-locked Ti-sapphire laser with an 80-MHz repetition rate (Tsunami, Spectra Physics), pumped by a 10-W Millennia (Spectra Physics) followed by a pulse-peaker (rate 8 MHz) and a third harmonic generator (Spectra-Physics, model 3980). The third harmonic beam was used for excitation of the sample inside the time-correlated single-photon counting instrument (IRF = 70 ps), and the second harmonic beam was collected as for the start pulse. Luminescence transients were fitted by a nonlinear least square fitting

procedure to a function
$$\left(X(t) = \int_{0}^{t} E(t')R(t-t')dt'\right)$$
 compris-

ing of the convolution of the IRF (E(t)) with a sum of exponentials

 $\left(R(t) = A + \sum_{i=1}^{N} B_i e^{-t/\tau_i}\right)$ with preexponential factors (B_i), characteristic lifetimes (τ_i) and a background (A). Relative concentra-

tion in a multiexponential decay was finally expressed as $c_n =$

 $\left(B_n / \sum_{i=1}^N B_i\right) \times 100$. The quality of the curve fitting is evaluated

by reduced chi-square and residual data.

For the determination of hydrodynamic diameters of the RMs, dynamic light scattering (DLS) measurements were made with Nano-S Malvern instrument using a 4-mW He-Ne laser $(\lambda = 632.8 \text{ nm})$. The size of the RM water pool was estimated by subtracting twice the molecular size of the surfactants (AOT = 1.1 nmand CTAB = 2.2 nm) from the respective measured diameter of the RM (assuming no solvent penetration) (Luisi and Magid, 1986).

RESULTS AND DISCUSSIONS

Figure 1a shows the absorption spectra of Rf bound to RBP in buffer and AOT-RMs of different degrees of hydration. Isoalloxazine ring is known to be responsible for the light absorption and emission of the flavin chromophore in the near-UV and visible regions (Choi and McCormick, 1980; Mataga et al., 2000). Although Rf absorbs mostly approximately at 450 nm, it also shows distinct absorbance in the 300-nm region (Figure 1a). The complexation of Rf with RBP in buffer is found to quench both the fluorescence of Rf (excited both at 300 and 445 nm) and tryptophan (Trp) fluorescence (excited at 300 nm) of protein (Figure 1b) compared with the free Rf and RBP in buffer, respectively (data not shown). The observed quenching of Rf upon binding with RBP in buffer is a consequence of ultrafast ET to the flavin chromophore (Rf) in the excited electronic state from nearby tryptophan or tyrosine residues present in RBP (Mataga et al., 2000; Zhong and Zewail, 2001; Lukacs et al., 2012; Rakshit et al., 2012). Upon confining RBP spatially in the AOT-RM of $w_0 = 10$, we find that Rf retains its characteristic fluorescence intensity at 520 nm (excited both at 300 and 445 nm; Figure 1b and its inset), which is characteristic of free Rf in buffer. However, in the AOT-RM of $w_0 = 20$ and 40, Rf fluorescence is quenched (Figure 1b and square symbol of Figure 1c), which is only observed when it binds to RBP and appears to mark the associated ET from RBP to Rf. The binding of Rf to RBP can similarly be concluded from the emission intensity (at 350 nm) of tryptophan residue of RBP, but this requires more



Figure 1. (a) Absorption spectra of RBP–Rf complex in buffer and AOT-RMs of $w_0 = 10$, 20 and 40. The black dashed line represents free Rf in buffer. (b) The total emission spectrum of the RBP–Rf complex in buffer and AOT-RMs of $w_0 = 10$, 20 and 40 ($\lambda_{ex} = 300$ nm). The inset shows the emission spectrum of the RBP–Rf complex excited at 445 nm in buffer (pink dotted line) and AOT-RMs of $w_0 = 10$, 20 and 40. (c) Relative fluorescence intensity of the Trp of RBP and RBP–Rf complex and the Rf of RBP–Rf complex in buffer and AOT-RMs of $w_0 = 10$, 20 and 40 (circle, RBP, $\lambda_{ex} = 300$ nm; triangle, RBP–Rf with $\lambda_{ex} = 300$ nm; square, RBP–Rf with $\lambda_{ex} = 445$ nm).

concern. Upon confining RBP in AOT-RM, we find that its fluorescence intensity (λ_{em} = 350 nm) increases in RM with decreasing w_0 values (Figure 1c, circular symbols), which might result from a change in the protein conformation due to confinement. On the other hand, in case of RBP–Rf complex, the fluorescence intensity of tryptophan residue increases in the AOT-RM of w_0 = 10 followed by substantial decrease in RM of w_0 = 20 and 40 (Figure 1c, triangular symbol). This quenching in tryptophan fluorescence in RM of w_0 = 20 and 40 could be related to Rf-binding in the protein nanospace and is believed to be both static and dynamic quenching discussed later.

To get a better insight into the observed change in fluorescence intensity, time-resolved fluorescence measurements are performed. Figure 2a shows the quenched decay transients of RBP–Rf complex (excited at 300 nm) in buffer compared with that of free RBP. The significant quenching of the Trp fluorescence is very much in line with the observed steady-state emission data discussed earlier (Figure 1b). Previously, Choi and McCormick (1980) suggested that the quenching of RBP fluorescence upon binding of Rf is mainly due to the ground-state stacking interaction between a Trp residue at the



Figure 2. (a) Fluorescence decay transients of tryptophan of RBP–Rf complex (λ_{ex} = 300 nm; decay monitored at 350 nm) in buffer (square) and AOT-RMs of w_0 = 10 (triangle) and w_0 = 40 (circle). The pentagon represents free RBP decay in buffer. (b) Fluorescence decay transients of Rf of RBP–Rf complex (λ_{ex} = 445 nm; decay monitored at 520 nm) in buffer (triangle) and AOT-RMs of w_0 = 10 (square), w_0 = 40 (circle).

binding site and the quinoxaline portion. In fact, the distinct shift in the absorption peak position of RBP-bound Rf, compared with free Rf in buffer, supports such an explanation (Figure 1a). However, the significant quenching of Trp fluorescence in the excited state (Table 1) suggests that the guenching is dynamic in nature also. It is to be noted that considerable geometrical confinement of RBP (without Rf) in RM of $w_0 = 10$ leads to an increase in the average lifetime ($<\tau>=a_1\tau_1+a_2\tau_2+a_3\tau_3$) of Trp residue (Table 1). However, increase in the water content of the RM reduces $<\tau>$ of Trp to reach a value comparable with that obtained in buffer (Table 1). On the other hand, when RBP-Rf complex is confined in RM, the Trp lifetime in RM of $w_0 = 10$ system is found to be more or less comparable with that of RBP itself confined in RM of $w_0 = 10$, which clearly identifies the inability of RBP to bind Rf at this hydration level of RM. However, it is evident from Table 1 that for RBP-Rf complex in RM of $w_0 = 20$ and 40, $<\tau >$ of Trp is guenched compared with that of RBP in RM with identical hydration, indicating that RBP is able to bind Rf in RM of $w_0 = 20$ and onwards.

The binding of Rf at a higher hydration of the protein is also evident in the absorbance (Figure 1a) and visible CD spectra (Figure 3c) of RBP–Rf complex in RM. In Figure 1a, it is observed that the bathochromic shift (~10 nm) in the visible band and a shoulder at ~490 nm, which is characteristic of Rf binding to RBP (Choi and McCormick, 1980), is absent in $w_0 = 10$ RM and is present in RM of $w_0 \ge 20$, indicating that RBP is able to bind Rf only in RM of $w_0 = 20$ and 40. Similarly, the appearance of strong CD bands in the visible region for RBP–Rf complex in RM of $w_0 = 20$ and 3d) suggests Rf binding to RBP.

$\lambda_{\rm ex}$ = 300 nm,	$\lambda_{\rm em}$ = 350 nm		<i>a</i> ₁	$ au_1$ (ns)	a2	$ au_2$ (ns)	<i>a</i> ₃	$ au_{3}$ (ns)	$<\!\!\tau\!\!>$ (ns
Buffer	_	RBP	0.49	0.24	0.43	1.18	0.08	3.64	0.92
	—	RBP-Rf	0.71	0.08	0.23	0.76	0.06	3.30	0.43
AOT-RM	$w_0 = 10$	RBP	0.33	0.24	0.54	1.50	0.13	4.44	1.47
	$w_0 = 20$		0.27	0.26	0.58	1.39	0.15	3.98	1.47
	$w_0 = 40$		0.42	0.16	0.47	1.35	0.11	4.23	1.17
	$w_0 = 10$	RBP-Rf	0.37	0.23	0.53	1.53	0.10	4.60	1.35
	$w_0 = 20$		0.37	0.21	0.54	1.42	0.09	4.21	1.22
	$w_0 = 40$		0.58	0.08	0.36	1.34	0.06	4.34	0.79
CTAB-RM	$w_0 = 25$	RBP	0.58	0.23	0.28	0.84	0.14	3.29	0.83
	$w_0 = 33$		0.58	0.21	0.28	0.78	0.14	3.22	0.78
	$w_0 = 25$	RBP-Rf	0.62	0.14	0.24	0.76	0.14	3.24	0.73
	$w_0 = 33$		0.60	0.14	0.27	0.68	0.13	3.33	0.69

 τ_i represents decay time constant, a_i represents its relative contribution and $\langle \tau \rangle$ is the average lifetime.



Figure 3. (a) Far-UV, (b) near-UV and (c) visible CD spectra of RBP in buffer and RBP-Rf complex in buffer, and AOT-RMs of w₀ = 10, 20 and 40. (d) Optical rotation value at 445 nm of RBP–Rf complex in buffer and AOT-RMs of $w_0 = 10$, 20 and 40. The inset shows the optical rotation value at 293 nm of RBP (square) and RBP-Rf (circle) in buffer and AOT-RMs of $w_0 = 10$, 20 and 40.

Importantly, the band positioned at ~445 nm is due to $\pi \rightarrow \pi^*$ transition, whereas those at 370 and 340 nm are attributed to a second $\pi \rightarrow \pi^*$ and $n \rightarrow \pi^*$ transitions, respectively (Figure 3c) (Galat, 1988). These strong CD bands suggest that Rf is rigidly packed in the binding cleft, and the rotation of the ribose moiety is completely hindered. It is important to mention here that the encapsulation of either RBP or RBP-Rf complex does not result in any significant change in the secondary structure (Figures 4a and 3a, respectively). However, the tertiary structure of the RBP-Rf complex as well as that of the protein itself suffers considerable perturbation upon confinement in $w_0 = 10$ RM system, as concluded from a loss in the intensity of the peaks at 268 and 293 nm (Figure 3b). With increased hydration ($w_0 \ge 20$), the near-UV CD spectrum considerably recovers its native form for the RBP-Rf complex, however, not for the protein itself (Figures 3b and inset of 3d). Indeed, one can observe that although the intensity of the CD signal for RBP at 293 nm is almost unaltered, it increases for the RBP-Rf complex with the increased hydration of the RM (Figure 3d, inset). This behaviour affirms that vitamin when bound to the protein indeed plays an important role in providing extra stability to protein in confined environment just like Rf-bound RBP has enhanced thermal stability manifested by the increase of denaturation temperature from 60.8°C to 72.8°C (Wasylewski, 2000).



Figure 4. (a) Far-UV CD spectra of RBP in buffer and AOT-RMs of $w_0 = 10$, 20 and 40. (b) Representative DLS signals for AOT and CTAB-RMs.

The observations mentioned earlier might be better understood in terms of confinement and water properties in RM. Water present inside the RM can broadly be distinguished as bound type (water molecules hydrogen bonded to the interface) and bulk type (Verma *et al.*, 2010; Saha *et al.*, 2012). At very low w_0 , the water molecules in the pool remain very close and strongly attracted to the polar head groups of the surfactants and hence are rather slow moving. As w_0 increases, the size of the water pool increases, and in such large water pools, the mobility of water molecules becomes relatively high. This change in mobility of the entrapped water molecules inside the RM has also been reported to affect the reaction kinetics occurring in the RM (Verma *et al.*, 2009). In RM of $w_0 = 10$, RBP retains its secondary structure although loses its tertiary structure partially as well as its binding capability. It is important to note here that the presence of hydration shell is essential for a macromolecule's biological activity (Pal and Zewail, 2004). Without hydration water, proteins would lack not only their native folded structure but also the conformational flexibility that allows their biological activity (Rupley and Careri, 1991). For proteins with nearly spherical shape, the empirical relationship between its molar mass (M) and the hydration degree of RM is described by the following relation (Eryomin and Metelitza, 1999),

$$w_0 = (0.083 \pm 0.008)\sqrt{M} \tag{1}$$

Equation (1) indicates that the maximum degree of hydration for RBP (M = 30 kDa) is expected to occur at $w_0 = 13-16$. It seems that in $w_0 = 10$ AOT-RM, RBP is in a lower degree of hydration compared with that in aqueous buffer, making the protein relatively rigid and hence not functional. In RMs of $w_0 \ge 20$, the abundance of bulk type water compensates its dehydration and also geometrical confinement is relieved at a higher extent compared with that in $w_0 = 10$ RM, bringing in RBP's native tertiary structure and hence its functionality.

The complexation of Rf with RBP in buffer leads to the quenching of Rf fluorescence (Mataga et al., 2002), as also observed in Figure 1b. Although significant quenching in the steady-state fluorescence of Rf is observed upon binding with RBP in $w_0 = 20$ and 40 RM, its lifetime does not change appreciably (Figure 2b and Table 2). The unchanged fluorescence decay of Rf in RBP rules out the occurrence of ultrafast ET in AOT-RMs ($w_0 \ge 20$). As has been discussed earlier, RBP regains its tertiary structure and binding capacity in RMs of $w_0 \ge 20$. So quenching in the lifetime of Rf bound to RBP is expected in RM of $w_0 = 20$ and 40 (Figure 2b). Remarkably, ET and its rate are dependent on the redox centre distance (Isied et al., 1992; Winkler, 2006; Mondol et al., 2012), for example, tunnelling times range from a few nanoseconds (12.2 Å ET in the high-potential iron-sulfur protein from Chromatium vinosum) to 10 ms (26 Å ET in Pseudomonas aeruginosa azurin) (Winkler, 2006). For flavin chromophores with various flavoproteins, the ET rate is found to be different due to different chromophorearomatic amino acid residue arrangements in the protein nanospace (Mataga et al., 2000). The donor-acceptor distance dependence of photoinduced ET in flavoproteins has been revealed by the work of Tanaka et al. (2007). In the confinement of AOT-RM ($w_0 \ge 20$), the nonoccurrence of ultrafast ET in the protein nanospace is thus

$\lambda_{\rm ex}$ = 445 nm,	$\lambda_{\rm em}$ = 520 nm		<i>a</i> ₁	$ au_1$ (ns)	<i>a</i> ₂	$ au_2$ (ns)	<i>a</i> ₃	$ au_3$ (ns)	$< \tau >$ (ns)
Buffer	_	Rf	0.00	1.00	0.00	1.00	1.00	4.70	4.70
	_	RBP-Rf	0.83	0.03	0.05	0.80	0.12	4.78	0.64
AOT-RM	$w_0 = 10$	Rf	0.20	1.11	0.00	1.00	0.80	4.84	4.09
	$w_0 = 20$		0.20	0.62	0.00	1.00	0.80	4.69	3.88
	$w_0 = 40$		0.15	0.79	0.00	1.00	0.85	4.72	4.13
	$w_0 = 10$	RBP-Rf	0.21	0.61	0.00	1.00	0.79	4.81	3.93
	$w_0 = 20$		0.20	0.65	0.00	1.00	0.80	4.81	3.98
	$w_0 = 40$		0.18	1.20	0.00	1.00	0.82	4.81	4.16
CTAB-RM	$w_0 = 25$	Rf	0.68	0.21	0.32	0.60	0.00	1.00	0.33
	$w_0 = 33$		0.76	0.27	0.24	0.81	0.00	1.00	0.40
	$w_0 = 25$	RBP-Rf	0.74	0.05	0.24	0.45	0.02	2.45	0.20
	$w_0 = 33$		0.73	0.05	0.24	0.52	0.03	2.44	0.23

 τ_i represents decay time constant, a_i represents its relative contribution and $<\tau>$ is the average lifetime.

believed to be the improper distance between the donor and the acceptor pair because of the electrostatic charge interaction of the RM surface as discussed later. Remarkably, in a proper redox distance with aromatic amino acids in RBP, the near-UV absorption band position of Rf does not undergo any shift (Choi and McCormick, 1980). The combination of red shift by complexing with aromatic amino acids at the binding site and blue shift due to nonpolar environment in the protein nanospace results in nearly unshifted band position of Rf in RBP. However, in AOT-RMs of $w_0 \ge 20$, a small blue shift in the near-UV band position coupled with a shoulder at 490 nm (Figure 1a) suggests the improper complexation of Rf with the aromatic amino acids. In this context, the steady-state quenching of Trp fluorescence in AOT-RMs also needs some attention. Previously, the observed substantial quenching of Trp fluorescence (~65%, data not shown) in buffer was attributed to

the presence of 5 out of 6 tryptophans in RBP in the vicinity (<0.45 nm) of Rf (Choi and McCormick, 1980). However, for AOT-RMs of $w_0 \ge 20$, such quenching of Trp fluorescence is found to be only ~40% (data not shown) compared with the free RBP in respective RMs. Thus, both steady-state absorption and emission spectroscopic data suggest the improper complexation of Rf for ET in the RBP interior, and in this condition, the observed steady-state quenching of Rf fluorescence in RM of $w_0 = 20$ and 40 systems is believed to be due to the ground-state stacking interaction of Rf with RBP.

To investigate the effect of confinement by a cationic surfactant forming RM, we chose CTAB-RMs with similar water pool size ($w_0 = 25$ and 33) to that of the AOT-RM of $w_0 = 10$ and 40, respectively (Table 3). Figure 4b is the representative DLS signals for each AOT and CTAB-RMs. The DLS signals are indicative of

Table 3. DLS data of AOT and CTAB-RMs of various w_0 values								
AOT-RM			CTAB-RM					
W ₀	Droplet diameter (nm)	Water pool size (nm)	w _o	Droplet diameter (nm)	Water pool size (nm)			
5	6.70	4.50	10	6.10	1.90			
10	7.56	5.36	20	8.40	4.20			
20	10.06	7.86	25	9.90	5.70			
40	14.90	12.70	30	14.4	10.2			
_	_	_	35	22.2	18.0			
_	_		40	24.0	20.3			



Figure 5. (a) Absorption spectra of RBP–Rf complex in CTAB-RMs of $w_0 = 25$ and 33. (b) Emission spectra of free Rf and RBP–Rf complex in CTAB-RMs of various w_0 values ($\lambda_{ex} = 445$ nm). (c) Visible CD spectra of Rf and RBP–Rf complex in buffer, and CTAB-RMs of $w_0 = 25$ and 33. (d) Fluorescence decay transients of tryptophan of RBP–Rf complex ($\lambda_{ex} = 300$; decay monitored at 350 nm) in CTAB-RMs of $w_0 = 25$ (square) and 33 (circle). The inset shows decay transient of Rf of RBP–Rf complex ($\lambda_{ex} = 445$; decay monitored at 520 nm) in CTAB-RMs of $w_0 = 25$ (square) and 33 (circle).

monodispersed RMs, and the observed values are guite consistent with previous literature (Corbeil and Levinger, 2003; Mitra et al., 2008). The fluorescence decay transients of RBP fluorescence in the RBP-Rf complex at different hydration levels of CTAB-RM is shown in Figure 5d. The observed quenching in the Trp fluorescence is evident from Table 1 and essentially indicates successful binding of Rf to RBP following the minimum hydration criterion of RBP in a RM (Equation (1)). In fact, the binding of Rf to RBP at these degrees of hydration is very much consistent with the associated absorption and CD spectroscopic data. In Figure 5a, the shoulder at ~490 nm, which is characteristic of Rf binding to RBP, is distinctly observed in both the hydration levels of CTAB-RM. Similarly, the strong CD bands at 445, 370 and 340 nm (Figure 5c) in the visible region for the RBP-Rf complex further supports the hindered rotation of the ribose moiety in the nanospace of RBP. In this situation, one can see the significant quenching of Rf fluorescence in both the steady-state and time-resolved fluorescence spectroscopic data of Figures 5b and 5d, respectively. The observed quenching of Rf fluorescence is well evident from the decay components of Table 2. Notably, the confinement of Rf in CTAB-RM itself guenches the Rf fluorescence (Table 2) and is found to be the guenching effect of bromide ion. In this respect, we monitor the decay transients of Rf in KBr solution and various RMs containing bromide ion, showing the guenching of Rf fluorescence compared with that in buffer (data not shown). It is to be noted that although Rf itself is guenched in the CTAB-RM, the decay time constants of free Rf in CTAB-RM are distinctly different than that present in the protein nanospace (Table 2). Hence, we emphasize that under the confinement of CTAB-RM, Rf present in the nanospace of RBP is free from the bromide guenching effect, and the observed quenching is essentially due to the ultrafast ET in the protein nanospace, which is similar to that in buffer (Table 2).

At this juncture, it is quite essential to discuss on the different behaviour of the flavoprotein under the confinement of RM with different charge types. One can see that the ET of Rf in CTAB-RM occurs in a manner similar to that in buffer, it certainly does not occur under the confinement of anionic AOT-RMs of similar hydration, although the hydration of the protein under confinement enables its binding to Rf. Recent simulation has shown that protein folding/unfolding equilibrium largely depends on the polarity of the confining volume (Tian and Garcia, 2011). The simulation by Griffin *et al.* (2005) suggests that in addition to the confinement effect, surface interaction plays determining role on the dimerization of an off-lattice β -barrel protein. In the present study, the different behaviour of the ET in various RMs suggests that the electrostatic interaction of the RM surface plays a significant role in determining the feasibility of such a process in flavoproteins. Notably in the RM pH ~7.0, RBP is expected to be negatively charged (*p*I of RBP is 4.0; Massolini *et al.*, 1995). A study by Lucent *et al.* (2007), using an explicit solvent model, showed that a protein is destabilized when confined by a purely repulsive potential together with solvent. Hence, we emphasize that in the repulsive potential of the confining volume (anionic AOT-RM), the binding site of the protein is so influenced that Rf is unable to complex properly with aromatic amino acids in the protein nanospace. However, under the attractive interaction of the cationic CTAB-RM, the protein conformation is so maintained that Rf can complex properly with aromatic amino acids in the protein nanospace, and ET is feasible with RBP.

CONCLUSION

It has been found that under the confinement of anionic AOT-RM, RBP loses its ability to bind with Rf in RM of $w_0 = 10$; however, it regains its binding capacity and tertiary structure in RM of w_0 \geq 20. The reason being that maximum degree of hydration for RBP occurs at $w_0 = 13-16$. Thus, when there is higher volume accessible to a protein and more bulk type water as in the AOT-RM of $w_0 = 20$ and onwards, RBP recovers its tertiary structure as well as binding capacity. Our detailed steady-state and timeresolved spectroscopic data suggest that under the similar size restriction of RM confinement, the biochemical function of RBP-Rf system is distinctly different depending on the hydration and the nature of the RM forming surfactant showing confinement. In anionic AOT-RM, confinement perturbs the ET from RBP to Rf, even at the maximum hydration of the protein. On the other hand, the cationic CTAB-RM resumes such a process with similar confining volume. We believe that because of the electrostatic repulsion of the anionic AOT-RM surface, the cofactor binding site of RBP is so perturbed that Rf is found unable to complex properly with the aromatic amino acids in protein interior. However, the confinement of cationic CTAB-RM resumes such a process with proper binding of Rf in the nanospace of RBP. Thus, both the hydration and surface charge of the confining volume is expected to largely determine the biochemical reaction dynamics like ET in real biological cells. Further studies will expand the present work to a wide range of biochemical reaction systems in different confinements of RMs.

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