MoS$_2$ Nanocrystals Confined in a DNA Matrix Exhibiting Energy Transfer

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3 Supporting Information

ABSTRACT: We report the wet chemical synthesis of MoS$_2$ nanocrystals (NCs), a transition-metal dichalcogenide, using DNA as a host matrix. As evidenced from transmission electron microscopy (TEM), the NCs are highly crystalline, with an average diameter of $\sim$5 nm. Ultraviolet-visible (UV-vis) absorption studies along with band gap calculations confirm that NCs are in quantum confinement. A prominent red shift of the optical absorption bands has been observed upon formation of the thin film using hexadecyltrimethylammonium chloride (CTAC), i.e., in the case of MoS$_2$@DNA–CTAC. In the thin film, strong electron–phonon coupling arises because of the resonance effect, which is reflected from the emergence of intense first-, second-, and third-order Raman peaks, whenever excited with the 488 nm line. We have established that our as-synthesized MoS$_2$ NCs quench the fluorescence of a well-known DNA minor groove binding probe, Hoechst 33258. Unprecedented fluorescence quenching (94%) of donor (Hoechst 33258) emission and efficient energy transfer (89%) between Hoechst 33258 and MoS$_2$ NCs (acceptor) are obtained. The donor–acceptor distance of these conjugates has been described by a Förster resonance energy transfer (FRET)-based model. Furthermore, employing a statistical method, we have estimated the probability of the distance distribution between the donor and acceptor. We believe that the study described herein may enable substantial advances in fields of optoelectronics, photovoltaics, catalysis, and many others.

■ INTRODUCTION

Transition-metal chalcogenide semiconductors display interesting properties that are important for sensing, catalysis, photovoltaics, and even biology.$^{1-8}$ In particular, molybdenum disulfide (MoS$_2$) is a prototypical transition-metal dichalcogenide material, which consists of covalently bonded Mo and S atoms. It is an indirect band gap semiconductor in its bulk form, with a band gap of 1.29 eV.$^9$ These nanomaterials have been known in the form of layered two-dimensional (2D) sheet, nested fulleren-like nanodots, inorganic nanotubes, or even as nanocrystals (NCs).$^{10-14}$ Although graphene is the best among all of the known 2D materials, MoS$_2$ nanosheets have proven to be certainly promising because they exhibit robust mechanical properties and superior electrical performance.$^{15,16}$ MoS$_2$ is also an effective lubricant because of its layered structure. Ultrathin layers of MoS$_2$ display strong photoluminescence that increases when the material is thinned from multilayer to monolayer because of the indirect–direct transition, arising from quantum confinement effects.$^{17}$ Nanostructured MoS$_2$ has been demonstrated to be an efficient catalyst for hydrogen evolution reaction (HER) and hydrodesulfurization (HDS).$^{18}$ The fact that all of these intense properties are combined in one material implies that MoS$_2$ could also be one of the most valuable materials in nanotechnology.

MoS$_2$ nanosheets, nanofibers, and nanorods have been prepared by a number of techniques, such as liquid exfoliation,$^{19}$ vapor deposition,$^{20}$ hydrothermal,$^{21}$ electrosprining,$^{22}$ wet chemical methods,$^{23}$ etc. At the same time, the colloidal synthesis of MoS$_2$ NCs is quite delicate, and as a result, there is less development compared to other synthetic routes. MoS$_2$ nanoparticles (NPs) having a size between 10 and 40 nm have been synthesized by the metal–organic chemical vapor deposition method.$^{24,25}$ Other approaches, such as sonochemical,$^{26}$ solvothermal,$^{27}$ or thermal decomposition methods,$^{28}$ do exist; however, the lack of size and shape control and produce poorly crystalline or totally amorphous products with little solvent dispersity. Yu et al. have reported less than 5 nm MoS$_2$ NCs, unfortunately dispersible only in organic solvents.$^{13}$ All of these methods involved complex or high-temperature reaction, and the morphology cannot be easily controlled. Therefore, it is still a challenge to develop a facile, room-temperature method to synthesize MoS$_2$ NCs.

The importance of finding a proper synthetic route for MoS$_2$ NCs has been a driving force for the present work. The use of biomolecules would be an interesting template for the synthesis of MoS$_2$ NCs because they are known to be nontoxic, viable.

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and excellent protecting agents. Quantum dots of ZnS, CdS, HgS, and PbS have been synthesized using various biomolecules. Among many biological systems that could participate in biomineralization and be indeed used as a biotemplate, DNA have been the subject of particular attention because of its excellent and predictable self-assembly properties, high rigidity of their double helices on the nanoscale, their stabilizing property, and most importantly, high affinity toward the metal cations. For example, self-assembled functionalization properties of DNA have enabled researchers to generate well-defined NPs, quantum dots, or even various superlattices. Herein, we report a facile wet chemical synthesis of MoS2 NCs using DNA macromolecules. To the best of our knowledge, this is the first time DNA has been used to synthesize and anchor MoS2 NCs. A high-resolution transmission electron microscopy (HRTEM) study reveals that NPs are highly crystalline in nature, with an average diameter of ∼5 nm. Ultraviolet–visible (UV–vis) absorption study and the corresponding band gap calculation demonstrate that as-synthesized NCs are in quantum confinement compared to the bulk MoS2. Moreover, our Raman spectroscopic studies indicate that the optically attractive MoS2 NCs could be Raman-active in strong resonance conditions. The presence of a strong Raman peak at 447 cm−1 along with second- and third-order Raman lines in the case of the MoS2 thin film further reinforces our justification on the resonance Raman effect. In the present study, we also establish that, akin to the other nanomaterials, MoS2 NCs could be an efficient fluorescence quencher. For that, we have chosen Hoechst 33258 (referred to as H258 henceforth), which is well-known as a potential DNA minor groove binder. It could be a Förster resonance energy transfer (FRET) donor when confined in DNA. Afterward, we show efficient energy transfer between H258 and proximal MoS2 NCs confined in DNA using steady-state and time-resolved spectroscopy. We first demonstrate a FRET-based model, which follows 1/r6 distance dependence, to calculate the donor–acceptor distance (r). We then explore the probability of donor–acceptor distance distribution using a simple mathematical model. Finally, employing the kinetic model developed by Tachiya (for the quenching of luminescent probes), we analyze the picosecond-resolved fluorescence results to understand the kinetics of energy transfer.
Salmon sperm DNA, sodium sulfide, sodium hydroxide, and MoCl5 were obtained from Sigma-Aldrich. Hexadecyltrimethylammonium chloride (CTAC) and H258 were obtained from Invitrogen. Dimethylandiline, hydrazine, and hydrochloric acid (HCl) were from Merck. Millipore water was used throughout the experiments. All of the chemicals were used as received without further purification. The H258–DNA solution was prepared by adding a requisite amount of the probe in DNA solution and stirring for 1 h. The final concentration ratio of [DNA]/[H258] was 10:1.

Synthesis of MoS2 NCs. It involves two steps. In the first step, an aqueous solution (5 mL) of DNA Na+ from salmon testes (65 mg) and 5 mL of 100 mM MoCl5 and H2O were mixed at room temperature under vigorous stirring, with a final pH value of ∼9 (adjusted by 1 M NaOH solution carefully). Then, the solution was allowed to stir for 6–8 h, so that the solution becomes colorless. In the second step, 5 mL of 200 mM Na2S was added to 10 mL of as-synthesized Mo−DNA complexes, with a final pH value of ∼6 (adjusted by HCl carefully), and the solution was stirred for 15 min. Completion of the reaction was observed visibly by color changes from colorless to yellow. Such a color transition is indicative of the spontaneous formation of the MoS2@DNA plate; and the solvent was evaporated slowly under the saturated vapor.

Preparation of the MoS2@DNA–CTAC Thin Film. An aqueous solution (5 mL) of as-prepared MoS2@DNA (final DNA concentration was 6.5 mg/mL) was added to 5 mL of aqueous solution of CTAC (6.5 mg/mL). A 1:1 stoichiometric combination led to the spontaneous formation of the MoS2@DNA–CTAC complex precipitate. The precipitate was collected by filtration, washed with distilled water, and then lyophilized for 48 h. Finally, a yellowish colorless solution was obtained when dried in a vacuum at 37 °C. The powder was then dissolved in n-butanol; the solution was cast on a Te plate; and the solvent was evaporated slowly under the saturated vapor at room temperature. Finally, a yellowish, water-insoluble thin film was formed, which was used for spectroscopic characterization.

Instrumentation. UV−vis Absorption Spectroscopy. Optical absorption spectra of the solutions were measured with a Shimadzu spectrophotometer using a quartz cuvette of 1 cm path length. Fluorescence Spectroscopy. Fluorescence spectra were recorded with Jobin Yvon Fluoromax-3 fluorimeter. Circular Dichroism (CD) Spectroscopy. The CD spectra were measured in a Jasco 815 spectropolarimeter. The cell path length was 10 mm.

Fourier Transform Infrared (FTIR) Spectroscopy. A JASCO FTIR-6300 spectrometer was used to confirm the covalent attachment of the DNA with the MoS2 NCs. For FTIR measurements, powdered samples were mixed with KBr powder and pelletized. The background correction was made using a reference of KBr pellets.

Raman Spectroscopy. Raman scattering measurements with the excitation laser lines of 488 and 785 nm were performed in a back-scattering geometry using a micro-Raman setup consisting of a spectrometer (model LabRAM HR, JobinYvon) and a Peltier-cooled charge-coupled device (CCD) detector. Raman spectra of all samples were recorded at room temperature in the frequency range of 100−4000 cm−1.

Transmission Electron Microscopy (TEM). TEM images were taken using a FEI TecnaiTF-20 field-emission high-resolution transmission electron microscope operating at 200 kV. Samples for TEM imaging were prepared by placing a drop of as-prepared aqueous MoS2@DNA solution on a carbon-coated Cu grid, and the solvent was evaporated under a light bulb.

Time-Correlated Single-Photon Counting. Picosecond-resolved fluorescence decay transients were measured using a commercially available spectrophotometer (Life Spec-ps, Edinburgh Instruments, U.K.) with 70 ps instrument response function (IRF). The observed fluorescence transients were fitted using a nonlinear least-squares fitting procedure to a function \( \chi(t) = J(E(t)R(t - t') d') \) comprising of convolution of the IRF \( E(t) \) with a sum of exponential \( R(t) = A + \sum_{i=1}^{N} B_i e^{-t'/\tau_i} \) with pre-exponential factors \( B_i \), characteristic lifetimes \( \tau_i \), and a background \( A \). The relative concentration in a multi-exponential decay was finally expressed as \( c_i = (B_i/\sum_{j=1}^{N} B_j) \times 100 \). The quality of the curve fitting was evaluated by reduced \( \chi^2 \) and residual data. It has to be noted that, with our time-resolved instrument, we can resolve at least one-fourth of the instrument response time constants after the deconvolution of the IRF.

RESULTS AND DISCUSSION

The process for synthesizing the MoS2 NCs is simple and involves two steps (see the Experimental Section for details). First, the addition of molybdenum chloride to the DNA solution followed by the increase of pH to 9 (by the addition of NaOH) under vigorous stirring renders the molybdenum ion to be coordinated with the DNA. After 6−8 h, the solution becomes colorless. Finally, sodium sulfide (Na2S) is added, and pH of the solution has been adjusted to 6 by adding HCl. The colorless solution turns yellow, indicating the formation of MoS2 NCs. We have analyzed the as-synthesized NCs by various microscopic and spectroscopic techniques. A typical TEM image of the MoS2 NCs is shown in Figure 1a. The NCs appear to be spherical in shape and fairly uniform in size. The particle sizes are estimated by fitting our experimental TEM data over 45 particles, which provides the mean diameter of ∼5 nm (Figure 1b). High-resolution (HR) images of a single particle are shown in panels c and d of Figure 1. The HRTEM images shown in panels f and g of Figure 1 as well as the selected area electron diffraction (SAED) demonstrate the crystalline nature of the as-synthesized particles. The distance between two adjacent planes is 0.25 nm, corresponding to the (102) lattice plane of hexagonal MoS2. Further confirmation regarding the composition of as-prepared NCs is also evident from energy-dispersive X-ray spectroscopy (EDAX). A X-ray diffraction (XRD) study (see Figure S3 of the Supporting Information) confirms the presence of hexagonal MoS2.

Figure 2a illustrates the UV−vis absorption spectra of DNA and MoS2@DNA in water. Well-defined absorption bands featuring at 384 and 468 nm appear for the final MoS2@DNA solution. Chikan et al. have reported a size-dependent spectroscopic study of MoS2 nanoclusters. They have established that MoS2 nanoclusters having a size between 3 and 8 nm have absorption maxima around 362−470 nm.

Figure 2. (a) UV−vis absorption spectra of DNA and MoS2@DNA in water. Well-defined absorption bands of the NCs are marked with arrows. (b) Band gaps from the onset of the absorption bands of the MoS2 NCs. (c) Photographs of DNA and MoS2@DNA solutions under visible light.
Because the two absorption peaks are present in the MoS2@DNA absorption spectrum, we associate the peaks with different sizes of MoS2 NCs rather than first- and second-order quantum confined states of the same NCs. This observation is also consistent with the broad size distribution obtained from our TEM studies. The optical band gaps of MoS2 NCs, measured from the onset of the absorption peaks, are ~2.30 and 2.51 eV, respectively (Figure 2b). Taking into account the band gap of bulk MoS2, i.e., 1.29 eV, one might expect photoluminescence from the as-prepared NCs because of the large quantum confinement effect. However, no such photoluminescence has been observed from our as-prepared NCs (data not shown), which reveals that evolution of photoluminescence because of the quantum confinement effect may not be applicable for our NCs; rather, it is more relevant for MoS2 nanosheets. To investigate MoS2 binding sites of DNA, we have performed FTIR studies. Figure 3 compares the FTIR of DNA and MoS2@DNA. Absorptions in the 1500–1250 cm⁻¹ region are caused by base-sugar vibrations. Sugar–phosphate vibrations appear in the 1250–1000 cm⁻¹ region. From the spectra, we can observe that both of the spectral regions have perturbed significantly, which indicates that there is a significant interaction of MoS2 with the DNA and that the samples are not merely mixtures of DNA and MoS2.

Raman spectra have been recorded at ambient temperatures using different excitation wavelengths. Figure 4 displays the Raman spectra of MoS2@DNA excited by 488 and 785 nm lines, respectively. It is clearly observed that there is a variation of peak intensity corresponding to the change in the excitation line. As depicted in Figure 4, the Raman spectra obtained upon excitation with the 488 nm line have peaks at 450 and 901 cm⁻¹, whereas a broad peak around 450 cm⁻¹ has been observed for the 785 nm excitation line. The evident peak broadening could be caused by the low spectral resolution of Raman spectroscopy with the 785 nm laser or because of the small size of the NCs integrated in large macromolecules. In contrast, in the case of 488 nm excitation, the Raman spectrum shows strong peaks, owing to the resonance Raman (RR) scattering, because the 488 nm line is in resonance with the band gap of the MoS2 NCs (2.3 eV). Furthermore, the second-order Raman peak at 901 cm⁻¹ also provides evidence about the RR effect.

Frey et al. disputed that the broad asymmetric peak around 450–460 cm⁻¹ for bulk MoS2 consists of two peaks, i.e., a second-order zone-edge phonon peak 2LA(M) at 454 cm⁻¹ and a first-order optical phonon peak A₂ at 465 cm⁻¹. Li et al. ascertained that the asymmetric feature splits into three peaks around 440, 450, and 459 cm⁻¹, where the 440 cm⁻¹ peak has been assigned to Mo=S vibrations for oxysulfide species and the later provided the supportive evidence for the argument by Frey et al. Most recently, Rao et al. have reported resonant Raman studies of a few layers of MoS2, where they have found that the 450–460 cm⁻¹ region becomes intense in resonance conditions. Taking into account the high-resonance conditions, we anticipate that the 450 cm⁻¹ peak observed for MoS2 NCs excited by the 488 nm line is due to the strong electron–phonon coupling, as expected in RR scattering.

Because 2D MoS2 has huge application in mechanics and field-effect transistors (FETs) at the nanoscale, we choose to prepare a thin film of MoS2 NCs. The procedure for the synthesis of the MoS2@DNA–CTAC thin film has been mentioned in the Experimental Section. The UV–vis absorption spectrum of the film has been shown in the inset of Figure 5. Because both CTAC and DNA have no UV–vis absorption after 300 nm implies that 395 and 484 nm peaks are solely due to the MoS2 NCs. Note that the absorption peaks have been red-shifted in the thin film, the effect of which is later reflected in the Raman spectrum of the MoS2@DNA–CTAC thin film. In comparison to solid MoS2@DNA, the MoS2@DNA–CTAC thin film has more intense Raman peaks (excited by the 488 nm line), as depicted in Figure 5. Raman spectra of DNA and CTAC alone have also been shown in Figure 5. While most of the Raman peaks of DNA and CTAC remain unaltered, a strong peak at 447 cm⁻¹ along with two other peaks at 896 and 1346 cm⁻¹ are clearly evident in the MoS2@DNA–CTAC thin film. The Raman peaks originating from the MoS2 thin film are much stronger as well as blue-shifted about 3–4 cm⁻¹ compared to MoS2@DNA in the solid state. In obvious contrast, the strong resonance condition arises in the MoS2 thin film, because the laser excitation energy coincides more with the electronic absorption band of the MoS2 thin film.
A very recent studies by Rao et al. demonstrated that MoS$_2$ could be a p-type conductor; therefore, it would prefer to interact with electron donor molecules.\textsuperscript{49} They have shown the charge-transfer interaction of MoS$_2$ with an electron donor molecule, such as tetrathiafulvalene (TTF), by monitoring the absorption band of TTF as well as MoS$_2$. We have also monitored the UV–vis spectrum of the MoS$_2$ NCs with various concentrations of a well-known electron donor; however, the absorption bands of MoS$_2$ NCs remain unaltered, indicating that MoS$_2$@DNA could not accept an electron in the ground state (see Figure S2 of the Supporting Information). It is widely accepted that nanomaterial confined in a biomolecule involves the energy transfer process, and several studies have been performed in this direction.\textsuperscript{48,50,51} Now, the question is, could MoS$_2$ NCs adequately quench the fluorescence of the dye/donor molecules, so that it can replace some of the currently used nanomaterials? In the present study, we have taken H258 as a donor, which is a well-known dye that interacts with the minor grooves of the DNA molecule. Steady-state fluorescence measurements have been carried out on both the H258–DNA and MoS$_2$@H258–DNA solutions. It has been found that fluorescence of H258 underwent drastic quenching in the presence of MoS$_2$ NCs. Figure 7b depicts the fluorescence spectra of H258 without and with the presence of MoS$_2$ NCs. As evidenced from the figure, an unprecedented fluorescence quenching of 94\% has been observed. The drastic quenching in the steady state obviously proves the efficacy of the MoS$_2$ NCs as a fluorescence quencher, however, does not provide any information about the nature of the quenching, i.e., whether it is static or dynamic. Picosecond-resolved fluorescence spectroscopy is a useful technique that provides information about the excited-state dynamics. Figure 7b shows the decay profiles of H258–DNA and MoS$_2$@H258–DNA monitored at the fluorescence maxima of the donor, i.e., at 470 nm ($\lambda_{\text{ex}} = 375$ nm). The decay transient of the donor (H258) has been fitted with two components, with an average lifetime of 2.6 ns. The fluorescence decay trace of the donor–acceptor (MoS$_2$@H258–DNA) could be fitted with a fast component, apparently corresponding to some non-radiative channel, along with two other components. The average lifetime of the donor–acceptor pair has been calculated to be 0.28 ns, which is much shorter than the average lifetime of the donor. Details of the fitting parameters of the time-resolved decays are tabulated in Table 1. Energy transfer, which involves deactivation of an electronic excited state of the donor and requires direct donor–acceptor spectral overlap, arises from the Columbic interaction between the donor and acceptor electric fields. Our study indeed demonstrates the incidence of a huge spectral overlap between the fluorescence of H258 (donor) and the absorbance of MoS$_2$ NCs (acceptor). The quenching of fluorescence decay transients and the spectral overlap, therefore, lead to the association of the energy transfer process.

The fitting result obtained from the decay transients is typical of a FRET donor decay transient, where rapid non-radiative transfer of energy from an excited donor to a ground-state acceptor manifests in a rapid drop off of the donor signal with time. We have estimated the energy transfer efficiency from the lifetime of the donor and donor–acceptor pair to be 89\%. The energy transfer efficiency obtained from the lifetime measurements is different from the steady-state measurements, however, more reliable, because the former is more sensitive than the latter as a result of the lamp fluctuation as well as many other processes. It is noteworthy to mention that we have compared to MoS$_2$@DNA, causing an enhancement in the total scattering cross-section.

To investigate any conformational and structural changes in DNA as a result of the formation of MoS$_2$ NCs, Circular Dichroism (CD) spectroscopy has been carried out. Figure 6 represents the CD spectra of DNA, in the absence and presence of MoS$_2$ NCs. As evidenced from the figure, a negative band at 246 nm and a positive band at 280 nm point toward the B form of DNA, consistent with the existing literature.\textsuperscript{45} Insignificant broadening with a slight red shift of the 280 nm band in the CD spectrum of MoS$_2$ indicates a little perturbation of the DNA structure. A decrease of molar ellipticity of the 248 nm band as well as a slight increase of 280 nm band intensity in the presence of MoS$_2$ NCs indicate the perturbation of the overall secondary structure of the DNA. The study of the DNA–NP interaction by Narayanan et al. and many others has also reported this kind of perturbation and argued that the perturbation may be associated with the condensation of the DNA in the presence of NPs because it could easily wrap around the NPs as a result of its flexible nature.\textsuperscript{46–48} Our study also provides supportive evidence for their argument.
calculated the donor−acceptor distance using the FRET method. Details of FRET have been described in the Experimental Section (see the Supporting Information). We have estimated the overlap integral $J(\lambda)$ to be $6.86 \times 10^{15} \text{ M}^{-1} \text{ cm}^{-1} \text{ nm}^{-4}$ according to eq 2 (see the Supporting Information). The characteristic Förster distance ($R_0$) is calculated to be 55 Å. Using the efficiency of FRET, we have calculated the donor−acceptor distance ($r$) to be 38 Å. Taking into account that the average radius of the MoS$_2$ NC is 2.5 nm (25 Å) (from the TEM measurements), the calculated donor−acceptor distance indicates that the residing probability of the donor H258 is very close to the surface of the NCs. To obtain an idea of the probability distribution of the donor−acceptor distance, we have analyzed the time-resolved decay transients of H258 in the presence and absence of MoS$_2$ NCs to construct the distance distribution function, $P(r)$ (see the Experimental Section for details). As evident in Figure 7d, the half width (hw) of the distance distribution is found to be 4 Å. This corresponds to a very high efficiency of energy transfer.

For a better understanding of the energy transfer between the excited state of H258 with MoS$_2$ NCs, it is essential to know the distribution of the acceptor (MoS$_2$ NCs) around the H258 molecules (bonded to the minor groove of DNA), because this is a governing factor that can influence the efficient energy transfer, as observed from the time-resolved fluorescence studies (Figure 7c). In this regard, we have applied a kinetic model developed by Tachiya for the quenching of fluorescent probes. $^{35,36}$ Figure 8 demonstrates the time-resolved fluorescence transients of H258 in the absence and presence of MoS$_2$ NCs, and black curves correspond to the fitting of the decay transients with eqs 15 and 16 (see the Experimental Section of the Supporting Information for details). The observed fluorescence transients were fitted using a nonlinear least-squares fitting procedure (software SCIENTIST) to a function $[X(t) = \int_0^t E(t')P(t - t')dt']$ comprising of the convolution of the instrument response function with the decay function.$^{35,36}$

### Table 1. Fitted Decay Time Constants of H258−DNA and MoS$_2$@H258−DNA from Picosecond-Resolved Experiments$^a$

<table>
<thead>
<tr>
<th>system</th>
<th>$\tau_1$ (ns)</th>
<th>$\tau_2$ (ns)</th>
<th>$\tau_3$ (ns)</th>
<th>$r_{av}$ (ns)</th>
</tr>
</thead>
<tbody>
<tr>
<td>H258−DNA</td>
<td>0.30 (25%)</td>
<td>3.40 (75%)</td>
<td>0.00</td>
<td>2.60</td>
</tr>
<tr>
<td>MoS$_2$@H258−DNA</td>
<td>0.02 (90%)</td>
<td>0.92 (6%)</td>
<td>6.1 (4%)</td>
<td>0.28</td>
</tr>
</tbody>
</table>

$^a$Values in parentheses represent the relative weight percentage of the time components. The standard error is $\sim$5%.
function (IRF) $[E(t)]$ with the exponential $p(t,m) = P(0)\exp\left\{-k_d - m_1[1 - \exp(-k_d t)] - m_2[1 - \exp(-k_d t)]\right\}$. The purpose of this fitting is to obtain the decays in an analytic form suitable for further data analysis. Reasonably good fitting has been observed from the model. The quenching parameters are summarized in Table 2. The quenching rate constant ($k_q$), which corresponds to unidentified traps, is same even after the addition of the acceptor (MoS$_2$ NCs), and this indicates the average number of unidentified traps around the donor H2S8 following a Poisson distribution. $k_q$ quenching rate constant by unidentified traps. $m$ is the mean number of H2S8–DNA molecules attached to one NC. $k_q$ is the rate constant for energy transfer for one H2S8–DNA molecule.

### Table 2. Overview of the Value of Quenching Parameters Using the Kinetic Model Developed by Tachiya

<table>
<thead>
<tr>
<th>System</th>
<th>$k_o$ (ns$^{-1}$)</th>
<th>$m$</th>
<th>$k_q$ (ns$^{-1}$)</th>
<th>$m$</th>
<th>$k_q$ (ns$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>H2S8–DNA</td>
<td>0.01</td>
<td>8.47</td>
<td>0.03</td>
<td></td>
<td>1.08</td>
</tr>
<tr>
<td>MoS$_2$@H2S8–DNA</td>
<td>0.01</td>
<td>8.47</td>
<td>0.03</td>
<td>1.08</td>
<td>6</td>
</tr>
</tbody>
</table>


“$k_o$ is the total decay constant of H2S8 in the excited state in the absence of the acceptor NCs. $m$ is the average number of unidentified traps around the donor H2S8 following a Poisson distribution. $k_q$ quenching rate constant by unidentified traps. $m$ is the mean number of H2S8–DNA molecules attached to one NC. $k_q$ is the rate constant for energy transfer for one H2S8–DNA molecule.”

## CONCLUSION

In conclusion, well-crystallized ~5 nm MoS$_2$ NCs have been synthesized in the DNA matrix. The calculated band gaps from the optical absorption of MoS$_2$ NCs confirm the quantum confinement of the NCs conjugated with the DNA matrix. A strong electron–phonon resonance condition makes MoS$_2$ NCs Raman-active. In comparison to MoS$_2$ NCs, the MoS$_2$ thin film has more intense Raman peaks because the laser excitation energy coincides more with the electronic absorption band of the MoS$_2$@DNA–CTAC thin film compared to MoS$_2$@DNA. Moreover, the present study also reveals that, similar to the other various nanomaterials, MoS$_2$ NCs would have a profound impact as an efficient fluorscence quencher. Analysis suggests that the fluorescence quenching of the donor in the presence of MoS$_2$ NCs is mainly due to a non-radiative decay channel, which confirms the energy transfer process. The donor–acceptor distance of 38 Å has been estimated using the efficiency of the FRET model. Further analysis of the probability of the donor–acceptor distance distribution suggests that the donor molecules are very close in proximity to the surface of the NCs. We have employed a kinetic model developed by Tachiya to understand the kinetics of energy transfer from H2S8 to MoS$_2$ NCs, assuming that the Poisson distribution of the quencher molecules around H2S8 bonded to the DNA minor grooves, which closely resembles the FRET data. Such a synthetic route of MoS$_2$ NCs as well as a thin film based on the DNA template may be extended to the other transition-metal chalcogenide materials, and we believe that MoS$_2$ NC-based energy transfer is expected to grow in the near future.

## ASSOCIATED CONTENT

## Supporting Information

Detailed experimental procedures, Raman spectra of MoO$_3$, and XRD of MoS$_2$@DNA. This material is available free of charge via the Internet at http://pubs.acs.org.

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### Notes

The authors declare no competing financial interest.

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